

The fluorescence emission spectrum of the europium (III) chelating polymer was recorded on a Shimadzu RF-5301PC spectrofluorophotometer.

3. Results and discussion

3.1 Fluorescence emission spectrum of europium (III) chelating polymer

The fluorescence emission spectrum of europium (III) chelating polymer presented in Fig. 2 was obtained under the excitation at 396 nm.

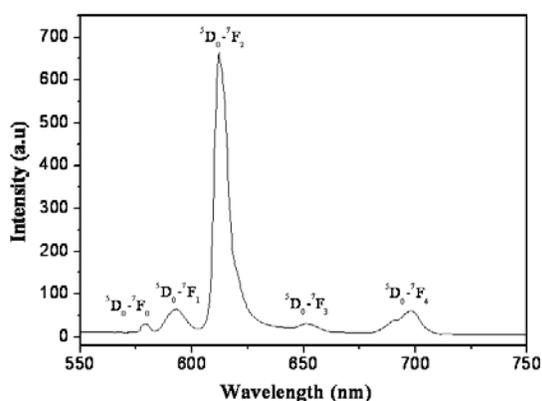


Fig. 2. Fluorescence emission spectrum of europium (III) chelating polymer.

As shown in Fig. 2, the ${}^5D_0 \rightarrow {}^7F_J$ transition with $J=0, 1, 2, 3$ and 4 are found at about 580 nm, 592 nm, 612 nm, 651 nm and 701 nm, respectively. The presence of only one ${}^5D_0 \rightarrow {}^7F_0$ line indicates that the Eu^{3+} ion occupies only a single site and a single chemical environment exists around it. The ${}^5D_0 \rightarrow {}^7F_1$ transition arises mainly from a magnetic dipole moment and is not strongly dependent on the site symmetry in which the Eu^{3+} ion is situated [7]. The ${}^5D_0 \rightarrow {}^7F_2$, ${}^5D_0 \rightarrow {}^7F_3$ and ${}^5D_0 \rightarrow {}^7F_4$ transitions are “forced” electric dipole transitions and may occur only at low symmetry sites [8]. The ${}^5D_0 \rightarrow {}^7F_2$ named “hypersensitive transition” by Jorgensen and Judd is especially sensitive to environment and increases in intensity with the amount of covalency between Eu^{3+} ion and the surrounding ligand and with decrease of site symmetry [8, 9]. The more intense ${}^5D_0 \rightarrow {}^7F_2$ transition of the europium (III) chelating polymer indicated that Eu^{3+} ion is in a site without a center of inversion.

3.2 Judd-Ofelt intensity parameters of europium (III) chelating polymer

According to Judd-Ofelt theory, the calculated oscillator strength for an induced electric dipole transition from the ground state to an excited state is [10]:

$$f = \frac{8\pi^2 m c \nu}{3h(2J+1)} \frac{(n^2+2)^2}{9n} \sum_{\lambda=2,4,6} \Omega_{\lambda} \left(\Psi J \parallel U^{\lambda} \parallel \Psi' J' \right)^2$$

where n is refractive index of the medium, J is the total angular momentum of the ground state, ν is the wave number of the transition, Ω_{λ} ($\lambda=2, 4, 6$) are Judd-Ofelt intensity parameters and $\parallel U^{\lambda} \parallel^2$ ($\lambda=2, 4, 6$) are the doubly reduced matrix elements evaluated in the intermediate coupling approximation for a transition $\Psi J \rightarrow \Psi' J'$. The transition intensity f depends on the $\parallel U^{\lambda} \parallel^2$ values between ΨJ and $\Psi' J'$ manifold. Due to selection rules and the unique nature of transition intensities for Eu^{3+} ion, any one of the $\parallel U^{\lambda} \parallel^2$ parameters decide the intensities of the transitions since the remaining two are zero. Thus, Ω_{λ} ($\lambda=2, 4, 6$) parameters can be evaluated independently from the emission transitions of ${}^5D_0 \rightarrow {}^7F_2$, ${}^5D_0 \rightarrow {}^7F_4$ and ${}^5D_0 \rightarrow {}^7F_6$, respectively [11]. The intensity of the ${}^5D_0 \rightarrow {}^7F_1$ fluorescent transition is found to be host independent whereas that of the ${}^5D_0 \rightarrow {}^7F_J$ transition depends on the Ω_{λ} parameters. This characteristic helps one to evaluate the Ω_{λ} parameters simply by the ratio of the intensities of the ${}^5D_0 \rightarrow {}^7F_{2,4,6}$ transitions to the intensity of ${}^5D_0 \rightarrow {}^7F_1$ transition as follows [10].

$$\int I_J d\nu / \int I_1 d\nu = A_J / A_1 = \frac{e^2}{S_{m d l}} \frac{\nu_J^3}{\nu_1^3} \frac{n(n^2+2)^2}{9n^3} \sum \Omega_{\lambda} \parallel U^{\lambda} \parallel^2$$

where $S_{m d l}$ refers to the strength of the magnetic dipole line strength of the ${}^5D_0 \rightarrow {}^7F_1$ transition in absorption.

The obtained values for parameter Ω_2 and Ω_4 of europium (III) chelating polymer are $13.9 \times 10^{-20} \text{ cm}^2$ and $4.733 \times 10^{-20} \text{ cm}^2$, respectively. Compared with other Eu^{3+} host-guest polymer system, europium (III) chelating polymer show a relative small Ω_2 value [12–14]. It has been well established that the Ω_2 parameter is structure sensitive and depends on the covalency of the rare earth ion sites [5]. It can be found in Fig.1 that the Eu^{3+} ion is coordinated by two oxygen atoms of the polymer ligands. While in Eu^{3+} host-guest polymer system, the Eu^{3+} ion is completely coordinated by the surrounding ligands atoms [12–14]. The decrease of Ω_2 value for europium (III) chelating polymer also indicates a decrease of the covalency between the Eu^{3+} ion and the surrounding ligands. The Ω_4 parameters have been related together to bulk properties of the lanthanide based hosts, but there is no theoretical prediction for this sensibility to macroscopic properties [15]. The Ω_6 intensity parameter was not determined because the ${}^5D_0 \rightarrow {}^7F_6$ transition could not be experimentally detected. This indicated that the Ω_6 is not important here.

3.3 Radiative transition probability of europium (III) chelating polymer

The electric (A_{ed}) and magnetic (A_{md}) dipole radiative transition probabilities can be evaluated from the following expressions [10]:

$$A_{ed} = \frac{64\pi^4 e^2 \nu^3}{3h(2J+1)} \frac{n(n^2+2)^2}{9} \sum_{\lambda=2,4,6} \Omega_{\lambda} (\Psi J \| U^{\lambda} \| \Psi' J')^2$$

$$A_{md} = \frac{64\pi^4 \nu^3}{3h(2J+1)} n^3 S_{md}$$

The sum of A_{ed} and A_{md} gives the radiative transition probability (A) for a transition $\Psi J \rightarrow \Psi' J'$ as:

$$A(\Psi J, \Psi' J') = A_{ed} + A_{md}$$

The total radiative transition probability (A_T) for an excited state is given as the sum of the $A(\Psi J, \Psi' J')$ terms calculated over all the terminal states.

$$A_T(\Psi J) = \sum_{\Psi' J'} A(\Psi J, \Psi' J')$$

As an excited state ΨJ is relaxed to several lower-lying states $\Psi' J'$, the radiative branching ratio β_R is defined as:

$$\beta_R(\Psi J, \Psi' J') = A(\Psi J, \Psi' J') / A_T(\Psi J)$$

The branching ratios can be used to predict the relative intensities of all emission lines originating from a given excited state. The experimental branching ratios can be found from the relative areas of the emission lines.

The rate of depopulation of an excited state is given by the radiative lifetime, $\tau_R(\Psi J)$:

$$\tau_R(\Psi J) = 1/A_T(\Psi J)$$

The stimulated emission cross-section, $\sigma(\lambda_p)(\Psi J, \Psi' J')$, between the states ΨJ and $\Psi' J'$ having a probability of $A(\Psi J, \Psi' J')$ is given by [10]:

$$\delta(\lambda_p)(\Psi J, \Psi' J') = \frac{\lambda_p^4}{8\pi c n^2 \Delta\lambda_{eff}} A(\Psi J, \Psi' J')$$

where λ_p is the wavelength of peak emission (in nm) and $\Delta\lambda_{eff}$ is the effective line width of the transition (in nm) found by dividing the integrated area of the emission band by its average height.

The predicted radiative transition probabilities of electric dipole transition A_{ed} and magnetic dipole transition A_{md} , the total transition probability $A_T(^5D_0)$, the fluorescence branching ratio β_R , the stimulated emission cross-section σ and the radiative lifetime τ_R are presented in Table 1.

Table 1. Radiative properties of europium (III) chelating polymer.

ΨJ	$\Psi' J'$	$\lambda(\text{nm})$	$\ U^2\ ^2$	$\ U^4\ ^2$	$\ U^6\ ^2$	$A_{md}(\text{s}^{-1})$	$A_{ed}(\text{s}^{-1})$	$A(\text{s}^{-1})$	$\beta(\%)$	$\sigma(\text{cm}^2)$	
5D_0	7F_0	580	0	0	0	0	0	0	0	0	
	7F_1	592	0	0	0	51.90	0	51.90	9.70	0.23×10^{-22}	
	7F_2	612	0.0032	0	0	0	417.06	417.06	77.93	2.71×10^{-21}	
	7F_3	651	0	0	0	0	0	0	0	0	
	7F_4	701	0	0.0023	0	0	66.23	66.23	12.37	1.98×10^{-22}	
	7F_5	-	0	0	0	0	-	-	-	-	
	7F_6	-	0	0	0.0003	0	-	-	-	-	
			$A_T(^5D_0) = 535.19 \text{ s}^{-1}$			$\tau_R = 1.868 \text{ ms}$					

* The $^5D_0 \rightarrow ^7F_5$ and $^5D_0 \rightarrow ^7F_6$ transitions could not be experimentally detected.

In Table 1, the transition $^5D_0 \rightarrow ^7F_2$ showed a β value of 77.93%. It has already established that an emission level with β value near 50% becomes a potential laser emission transition [16]. The most important parameter determining the potential laser performance at room temperature is the stimulated emission cross section σ . The large stimulated emission cross section is an attractive feature for low-threshold, high-gain applications [17]. The

stimulated emission cross section at the 612 nm for europium (III) chelating polymer is $2.71 \times 10^{-21} \text{ cm}^2$, which is comparable with those for Er^{3+} doped laser glasses [18].

5. Conclusion

In conclusion, a europium (III) chelating polymer has been synthesized and its radiative

properties has been studied. The Judd-Ofelt phenomenological parameters, Ω_2 and Ω_4 were obtained from the fluorescence emission spectrum are $13.9 \times 10^{-20} \text{ cm}^2$ and $4.733 \times 10^{-20} \text{ cm}^2$, respectively. Radiative properties of europium (III) chelating polymer were investigated also. The high fluorescence branching ratio of ${}^5\text{D}_0 \rightarrow {}^7\text{F}_2$ transition and large emission cross section showed that europium (III) chelating polymer is a promising material for use in optical devices.

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