

# Temperature-dependent electrical properties of Au Schottky contact and deep level defects in n-type GaN

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The Au/n-GaN Schottky barrier diode (SBDs) has been characterized by current-voltage (I-V), capacitance-voltage (C-V) and deep level transient spectroscopy (DLTS) measurements in the temperature range of 85 K - 405 K. The estimated Schottky barrier height is 0.25 eV at 85 K and 1.06 at 405 K by I-V method, respectively. Calculations showed that the barrier height is 1.02 eV at room temperature by C-V method. Little variation with either temperature or frequency was observed, and the carrier concentration derived from these measurements confirmed the dopant concentration of  $\sim 6.2 \times 10^{16} \text{ cm}^{-3}$ . A dominant trap  $E_a$ , with activation energy of 0.92 eV is observed in n-type GaN grown by metalorganic chemical vapor deposition (MOCVD). The capture cross section of  $1.9 \times 10^{-15} \text{ cm}^2$  is estimated based on the well-known logarithmic dependence of DLTS peak height with different filling pulse width for trap  $E_a$ . The trap  $E_a$  (0.92 eV), which is commonly observed in thin GaN layer grown by various techniques, is believed to be associated with nitrogen interstitial defect.

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## 1. Introduction

Due to superior characteristics of gallium nitride (GaN) such as direct wide band gap (3.4eV at room temperature), high break down field and high saturation velocity, it is one of the promising materials for the fabrication of optoelectronic and high temperature/high-power microelectronic devices like light emitting diodes (LEDs) [1], laser diodes (LDs) [2], photo detectors [3], metal oxide semiconductor field effect transistors (MOSFETs) [4], hetero junction field effect transistors (HFETs) [5], high electron mobility transistor (HEMTs) [6] and Schottky rectifiers [7]. However, the performance of GaN devices can be limited by the deep levels present in the GaN band gap, which act as either carrier trap or generation/recombination centers. Although there has been a substantial knowledge about the presence of these deep level defects in GaN [8, 9], it is essential to understand these defects to address issues related to the performance and reliability of GaN devices.

Many researchers have characterized and reported the defect levels in GaN which was grown by hydride vapor phase epitaxy (HVPE) [10] and metalorganic chemical vapor deposition (MOCVD) [11, 12, 13]. Gotz *et al* [11] reported the two deep levels with activation energies 0.49eV and 0.18eV with trap concentration  $6.3 \times 10^{14} \text{ cm}^{-3}$  for  $E_1$  and  $7 \times 10^{13} \text{ cm}^{-3}$  for  $E_2$ . Look *et al* [14] created a point defect in GaN by electron irradiation and reported electron traps at 0.18eV and 0.9eV by DLTS technique. He also reported that the electron traps at 0.18eV and 0.9eV could be related to nitrogen interstitial ( $N_i$ ) or gallium interstitial complex ( $Ga_i - X$ ). Goodman *et al* [15] identifies the nitrogen vacancy with activation energy of 0.20eV below the conduction band in n-GaN epitaxial

lateral over growth metallo-organ vapor phase epitaxy (ELO-MOVPE). Look *et al* [16] identify the trap having the activation energy  $90 \pm 2 \text{ meV}$  with effective capture cross section  $\sigma = 3 \pm 1 \times 10^{-22} \text{ cm}^2$  in GaN. Fang *et al* [17] observed a dominant trap with activation energy of 1.0 eV and apparent capture cross section  $2 \times 10^{-12} \text{ cm}^2$  in heterostructure AlGaIn/GaN/SiC. Py *et al* [18] reported the four known levels with activation energy in the range 0.1 - 0.94 eV in n-GaN which was grown by hydride vapor phase epitaxy (HVPE) by deep level transient spectroscopy. In this paper we report on the temperature-dependent electrical properties of Au Schottky contacts and characterize deep level defects in n-type GaN.

## 2. Experimental details

Gallium nitride (GaN) samples used in this study were grown by metalorganic chemical vapor deposition (MOCVD) on c-plane  $\text{Al}_2\text{O}_3$  substrate and were followed by the growth of  $2 \mu\text{m}$  n-type GaN with Si. The n-GaN layer was first ultrasonically degreased with warm trichloro ethylene followed by acetone and methanol for 5min each. This degreased layer was then dipped into boiling aquaregia [ $\text{HNO}_3 : \text{HCl} = 1:3$ ] for 10 min to remove the surface oxide and the sample was rinsed in deionized water. Ohmic contact Ti (20nm)/Al (40nm) was deposited on a portion of the sample by electron beam evaporation system under the vacuum of  $6 \times 10^{-6} \text{ mbar}$ . Then the sample was annealed at 750 °C for 2 min in nitrogen atmosphere. Schottky contacts are formed by the evaporation of gold (Au) as dots with a diameter of 1mm through stainless steel mask with a thickness of 60 nm. The current – voltage (I-V) characteristics were first measured in the

temperature range 85 to 405K by the step of 40K in the dark. Then, capacitance - voltage (C-V) and DLTS measurements were carried out by automated DLS-83D system (Semi Lab, Lock-in amplifier based system which facilitated measurement at pulse frequency in the mHz range).

### 3. Result and discussion

Typical current-voltage characteristics of gold (Au) Schottky contact on n-GaN in the temperature range of 85K - 405K are shown in Fig. 1. The current flow through the diode can be described by thermionic emission (TE) theory and I-V relationship of a Schottky diode is given by [19]

$$I = I_s \exp\left(\frac{qV}{nkT}\right) \left[1 - \exp\left(\frac{-qV}{kT}\right)\right] \quad (1)$$

where

$$I_s = AA^{**} T^2 \exp\left(\frac{-q\phi_{b0}}{kT}\right) \quad (2)$$

where,  $I_s$  is the saturation current,  $q$  is the electron charge,  $V$  is the applied voltage,  $T$  is the absolute temperature,  $n$  is ideality factor,  $k$  is Boltzmann's constant,  $A$  is the diode area,  $A^{**}$  is the effective Richardson constant ( $26.4 \text{ A cm}^{-2} \text{ K}^{-2}$  for n-GaN) [20] and  $\phi_{b0}$  is the Schottky barrier height (SBH). The value of  $\phi_{b0}$  can be deduced directly from the I-V curves, if the Richardson constant  $A^{**}$  is known. The I-V measurement were made to determine the saturation current  $I_s$  from which the barrier height was defined in terms of the thermionic emission theory

$$\phi_{b0} = \frac{kT}{q} \ln\left(\frac{AA^{**} T^2}{I_s}\right) \quad (3)$$

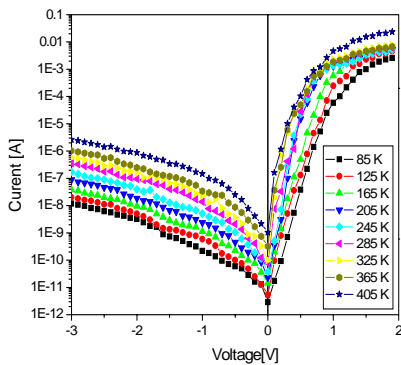


Fig. 1. Typical forward and reverse current-voltage (I-V) characteristics of Au Schottky contact on n-type GaN in the temperature range of 85-405 K.

Equation (1) shows that the logarithmic plot of  $I/[1 - \exp(-qV/kT)]$  against  $V$  as shown in Fig. 2 is a linear and  $I_s$  was obtained from the y-axis intercept at zero voltage. Using these  $I_s$  value the barrier height is determined from the equation (3). The calculated barrier height is 0.25 eV at 85K and 1.06 eV at 405K respectively. It is observed that the leakage current at -1V is  $2.18 \times 10^{-10} \text{ A}$  at 85K and  $1.47 \times 10^{-7} \text{ A}$  at 405K respectively. Many models have been proposed to explain the excess leakage current in Schottky diodes on GaN or AlGaIn/GaN, including the thin surface barrier model [21] and the field-emission/trap-assisted tunneling model [22]. However the detailed analysis of the I-V curves is beyond the scope of this paper.

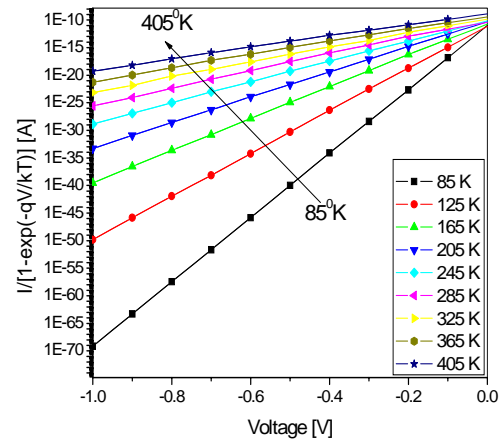


Fig. 2.  $\ln I/[1 - \exp(-qV/kT)]$  versus  $V$  plot for Au Schottky contact on n-GaN in the temperature range 85-405K.

The capacitance-voltage characteristics of Au Schottky contact on n-GaN were measured at temperatures ranging from 85 to 405 K. Little variation with either temperature or frequency was observed. Fig. 3 shows only one C-V curve which is measured at room temperature to check the quality of the Schottky diode. The barrier height was obtained by the following capacitance and voltage relation [20]

$$\frac{1}{C^2} = \left(\frac{2}{\epsilon_s q N A^2}\right) \left(V_{bi} - \frac{kT}{q} - V\right) \quad (4)$$

where  $\epsilon_s$  is the permittivity of the semiconductor ( $\epsilon_s = 9.5 \epsilon_0$ ) and  $V$  is the applied voltage. From the plot of  $1/C^2$  versus  $V$  gives the x-intercept  $V_0$  related to the built in potential by the relation  $V_{bi} = V_0 + kT/q$ , where  $T$  is the absolute temperature. The barrier height is given by the equation  $\Phi_b = V_0 + V_n + kT/q$ , where  $V_n = (kT/q) \ln(N_c/N_d)$ . The density of states in the conduction band edge is given by  $N_c = 2(2\pi m^* kT/h^2)^{3/2}$ , where  $m^* = 0.22m_0$  and its value is  $2.6 \times 10^{18} \text{ cm}^{-3}$  for GaN [23]. The calculated barrier height and carrier concentration of Au Schottky contact are 1.02eV and  $6.2 \times 10^{16} \text{ cm}^{-3}$ .

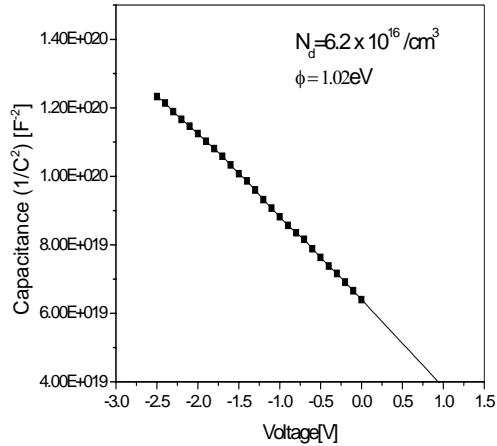


Fig. 3. The capacitance - voltage (C-V) characteristics of Au Schottky contact on n-GaN at room temperature.

To investigate whether there are electrically active as-induced defects generated in the band-gap, we performed deep level transient spectroscopy (DLTS) measurements. Any deep level can be experienced as either capture or emission of carrier, with transition probabilities per unit time depending on the free carrier concentrations, capture cross section of the trap, ionization energy and temperature. According to the value of applied voltage for the electron trap studied here in the space charge zone of Schottky diode on n-type GaN in the dark, either emission or capture of electrons toward or from the bottom  $E_C$  of the conduction band is relevant. When the Fermi level just coincides with the trap level  $E_T$ , the capture coefficient equals the emission rate  $e_n$  resulting at equilibrium which is given by [24]

$$e_n = \sigma_n v_{th} N_C \exp\left(\frac{E_T - E_C}{kT}\right) = \frac{1}{\tau_e} \quad (5)$$

here  $\sigma_n$  is the capture cross-section,  $v_{th}$  is the thermal velocity of an electron defined as  $v_{th} = (3kT/m_{eff})^{1/2}$ ,  $N_C$  is the effective density states at the bottom of the conduction band,  $k$  is the Boltzmann's constant and  $T$  is the absolute temperature,  $\tau_e$  is the emission time constant, which is the reciprocal of the emission time  $e_n$ . In order to characterize the deep level defects, DLTS measurement were performed over the temperature range 85-405K using automated DLTS system. Typically, a quiescent reverse bias of  $V_r = -1V$  was employed with filling -pulse voltage  $V_p = 0V$  and filling pulse width  $t_p = 0.1ms$ . Fig. 4 shows DLTS spectrum measured with lock-in frequency of 120 Hz. Only one prominent deep level in the as-grown GaN that peaked at 361 K is observed and labeled as ' $E_a$ ' is shown in figure 4. The basic DLTS analysis relies on spectrum which is obtained from the plot of the possible correlation of the emission transient (formerly box-car or lock-in) as either a function of the temperature [25] or the rate window [26], the other parameters being constant.

From several such spectra, the product  $e_n^{-1} v_{th} N_C$  can be calculated and plotted as a function of  $1000/T$ , leading usually to a straight line an Arrhenius plot. Fig. 5 shows an Arrhenius plot for one deep level (at 361K). It can be seen from figure 5 that the Arrhenius analysis the activation energy for the trap  $E_a$  was determined to be 0.92eV. We interpret the  $E_a$  as corresponding to the defect level with activation energy of 0.95eV which is similar to that reported by Goodman *et al* [27] and 0.879eV reported by Lee *et al* [28]. Fang *et al* [29] reported that the electron trap might be a  $N_I$  related defect created by the reaction  $N_N V_N + N_I$  [30] having the activation energy of 0.85eV. Auret *et al* [31] also reported the electron trap with activation energy of 0.92 eV with capture cross section of  $7.9 \times 10^{-14} \text{ cm}^2$  which is similar to the estimated activation energy (0.92 eV) of our samples.

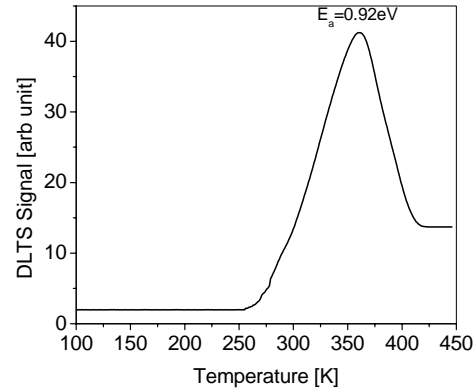


Fig. 4. DLTS spectrum of an Au Schottky contact on n-GaN grown by MOCVD. ( $V_r = -0.5V$ ,  $V_p = 0V$ ,  $t_p = 100 \mu s$ ,  $f = 120Hz$ ).

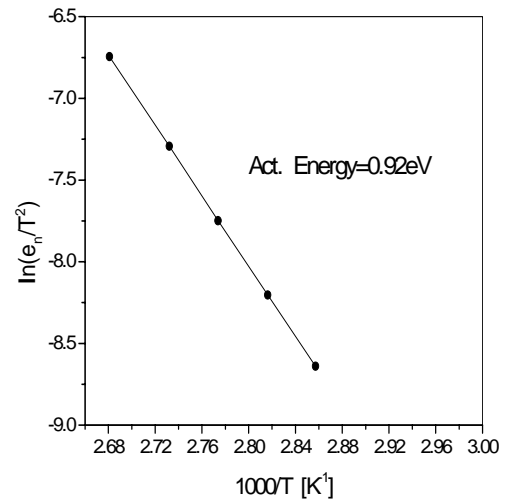


Fig. 5. The Arrhenius plot of  $\ln(e_n/T^2)$  versus  $1000/T$  of Au Schottky contact on n-GaN.

The carrier capture cross section of semiconductor deep level defect contains useful information on the nature of the defect. This parameter directly determined by measuring the capture time constant  $\tau_c$  of the deep level [32]

$$\tau_c = \frac{1}{\sigma_n v_{th} n} \quad (6)$$

where  $\sigma_n$  is the cross section,  $v_{th}$  is the thermal velocity of electron and  $n$  is the free carrier concentration at the measurement temperature.  $\tau_c$  is often measured by monitoring the deep level transient spectroscopy (DLTS) peak height as a function of filling pulse width  $t_p$  [33] (as shown in the figure 6), the slope of the plot  $\ln(1-N(t_p)/N_T)$  versus  $t_p$  as shown in figure 7 giving  $\tau_c^{-1}$  where  $N(t_p)/N_T$  is the peak amplitude. The above plot is usually found to be non linear due to different capture rate of the free carrier tail extending into the depletion region. The initial part of the plot, which is fairly linear, gives a reliable estimate of the  $\tau_c$  [34]. The capture cross section of  $E_a$  as determined from the figure 7 is  $1.9 \times 10^{-15} \text{ cm}^2$ . We interpreted the estimated capture cross section of our samples which were grown by MOCVD which is found to similar to the capture cross section ( $\sigma_a = 4.3 \times 10^{-15} \text{ cm}^2$ ) reported by Goodman *et al* [15] and reported by Py *et al* [18] ( $\sigma_a = 1 \times 10^{-16} \text{ cm}^2$ ).

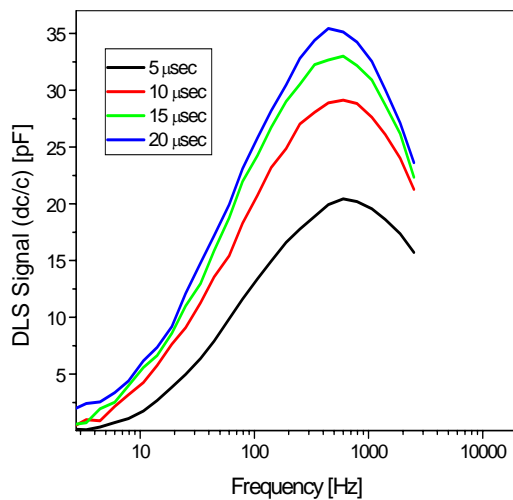


Fig. 6. Capture cross section measurement of DLTS spectra with different filling pulse width (5 $\mu$ s, 10 $\mu$ s, 15 $\mu$ s, 20 $\mu$ s,) graph of Au/n-GaN Schottky barrier diode.

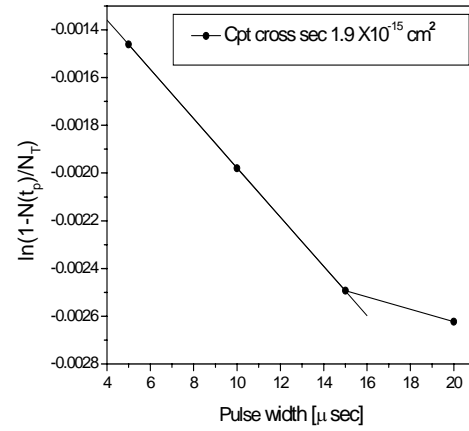


Fig. 7. The capture cross section plot of  $\ln(1-N(t_p)/N_T)$  versus Pulse width of Au/n-GaN Schottky diode.

The exact origin of the deep level remains an open question. Our observed level  $E_a$  is interpreted as corresponding to the 0.95 eV level measured by Goodman *et al* [27], and the 0.879eV level measured by Lee *et al* [28] and the 0.92 eV level measured by Auret *et al* [31]. They have suggested that the  $E_a$  level may be associated with a native defect in GaN and it has been assigned to nitrogen interstitial. In comparing DLTS measurements it would not be unexpected for electronic levels in GaN to grow by a variety of techniques. Conversely, the presence and concentration of defect levels associated with impurities might be expected to vary in metals grown by different techniques.

#### 4. Conclusions

We have investigated the electrical properties of Au/n-GaN Schottky diode in the temperature range of 85 K - 405 K. Little variation with either temperature or frequency was observed, and the carrier concentration derived from these measurements confirmed the dopant concentration of  $\sim 6.2 \times 10^{16} \text{ cm}^{-3}$ . Measurements showed that the barrier height of Au/n-GaN Schottky diode is 0.25 eV at 85 K and 1.06 eV at 405 K by I-V method, and 1.02 eV at room temperature by C-V method. Based on temperature-dependent I-V characteristics, it was observed that the barrier height is increased with increase in temperature. Deep level transient spectroscopy was used to characterize electronic defects in n-type GaN. A dominating trap  $E_a$  with activation energy of 0.92eV, which is commonly observed in thin GaN layer grown by various techniques, has been found in n-type GaN grown by metalorganic chemical vapor deposition. The capture cross section  $1.9 \times 10^{-15} \text{ cm}^2$  was extracted based on the logarithmic dependence of DLTS signal with filling pulse width for the trap  $E_a$ . The trap  $E_a$  is believed to be associated with nitrogen interstitial native defect. Level  $E_a$

is interpreted as corresponding to electronic state previously observed in GaN grown by HVPE and MOCVD.

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